

Topological features in photonics



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Fai della Paganella
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- Sunil Mittal, Vikram Orre, M. Khan, H. Kim, B. Cao
- S. Ganeshan, J. Fan, P. Adhikari, J. Taylor, A. Migdall (NIST/JQI)
- Vaezi (Cornell)
- M. Lukin, E. Demler (Harvard)
- P. Rabl (Vienna), C. Galland (Stuttgart)
- N. Harris, D. Englund (MIT)
- E. Kapit (CUNY), S. Simon (Oxford)
- F. Grusdt, F. Letscher, M. Fleischhauer (Kaiserslautern)



[publications: hafezi.umd.edu](http://publications:hafezi.umd.edu)

- Synthetic gauge field, formalism
- Experiment in silicon photonics
- Interacting regime: Fractional quantum Hall
 - Laughlin, Pfaffian etc.
 - Preparation
- Measuring topological invariants

Review: Integer Quantum Hall effect

Is it possible to have similar effects for photons?

Microwave : Haldane, Raghu (2008), Soljacic (2009)

Optical domain: weak magneto-optic effect!

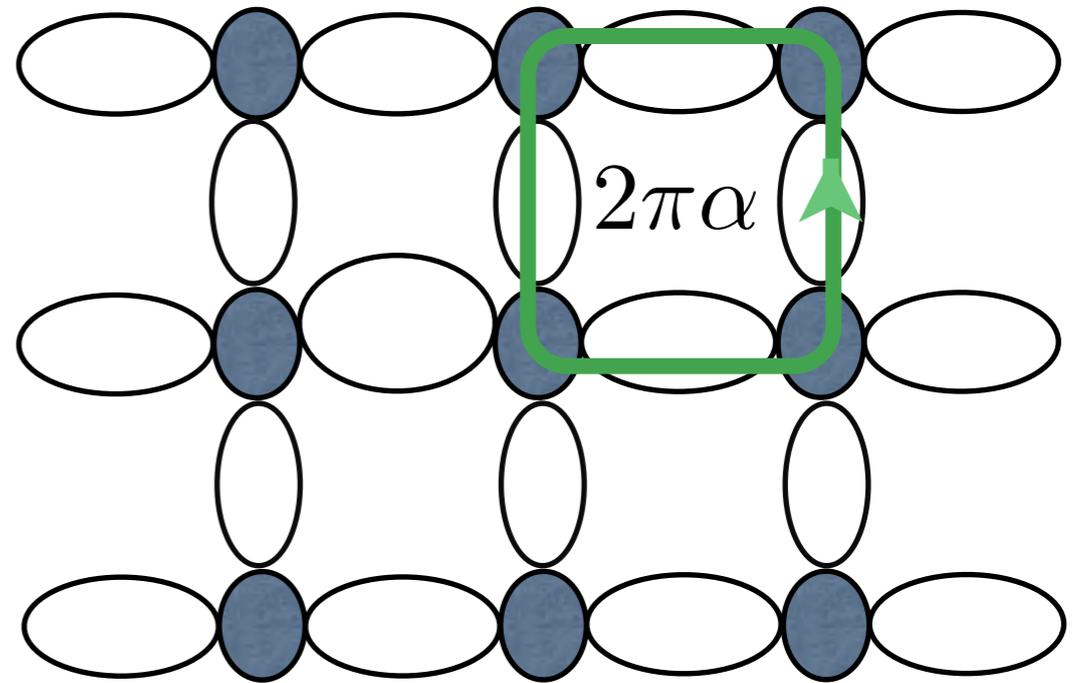
Standard for electrical conduction

Synthetic Magnetic Field

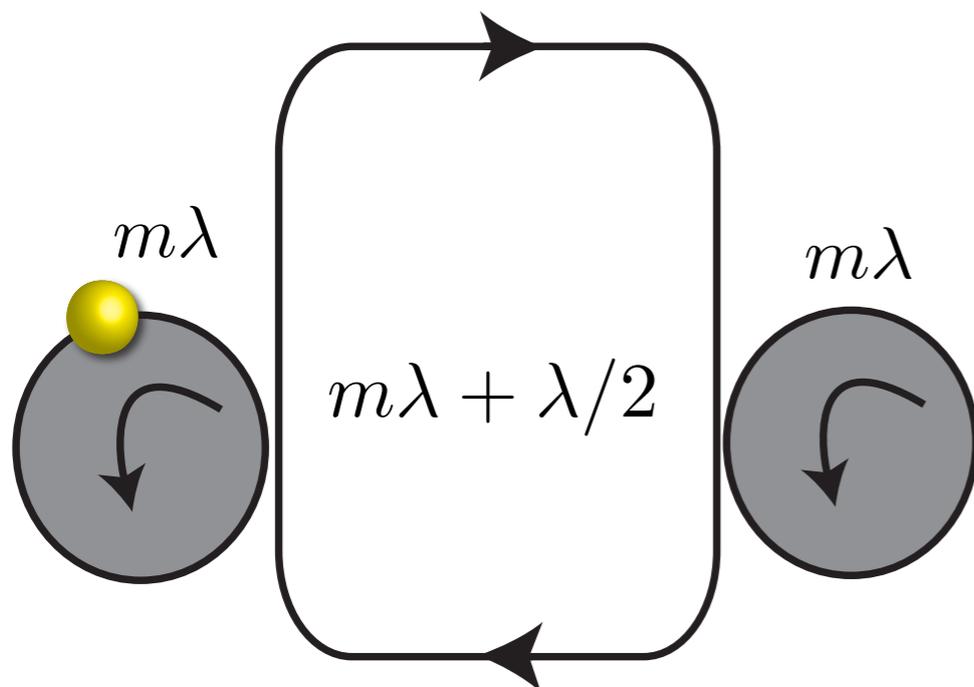
In analogy to electrons on a magnetic lattice:

$$H_0 = -J \sum_{x,y} \hat{a}_{x+1,y}^\dagger \hat{a}_{x,y} e^{-i2\pi\alpha y} + \hat{a}_{x,y}^\dagger \hat{a}_{x+1,y} e^{i2\pi\alpha y} + \hat{a}_{x,y+1}^\dagger \hat{a}_{x,y} + \hat{a}_{x,y}^\dagger \hat{a}_{x,y+1}$$

- Tight-binding form
- Magnetic phase

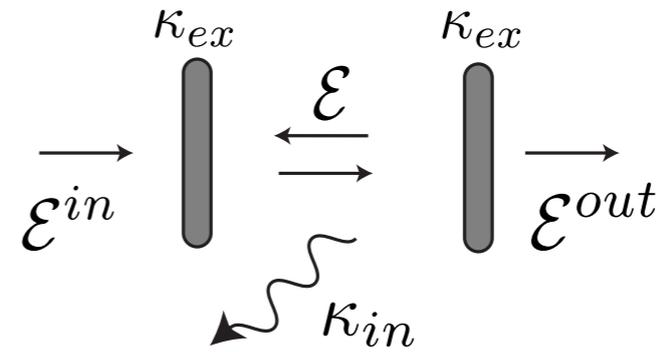
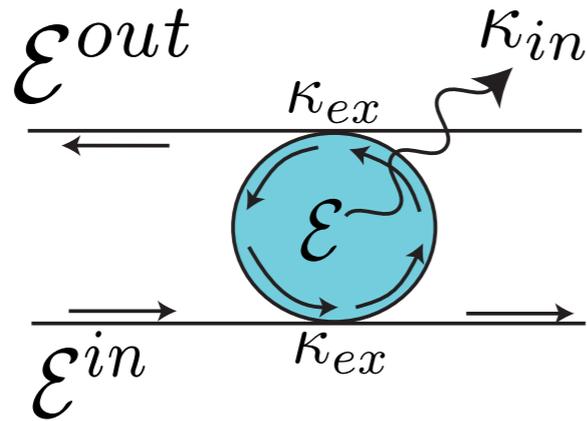


Two resonator case:



$$H_{eff} = -\kappa \hat{a}_{x+1}^\dagger \hat{a}_x e^{-2\pi i \alpha} - \kappa \hat{a}_x^\dagger \hat{a}_{x+1} e^{2\pi i \alpha}$$

Single-mode resonators: input-output formalism

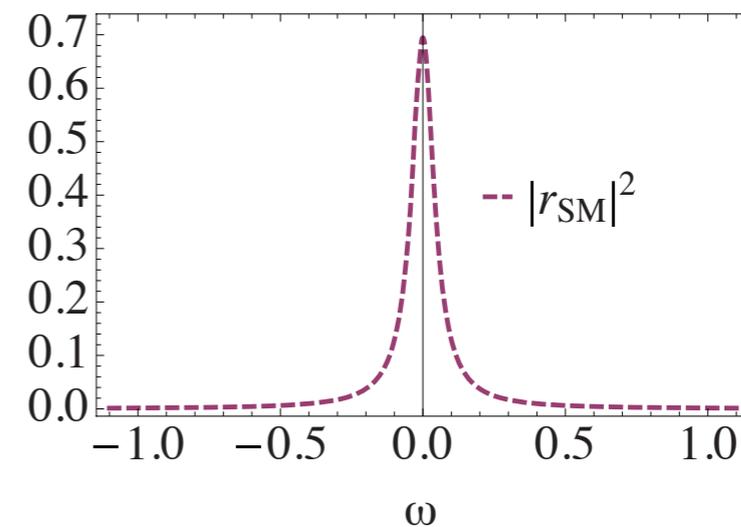


$$\frac{d\mathcal{E}}{dt} = (-\kappa_{in} - 2\kappa_{ex})\mathcal{E}$$

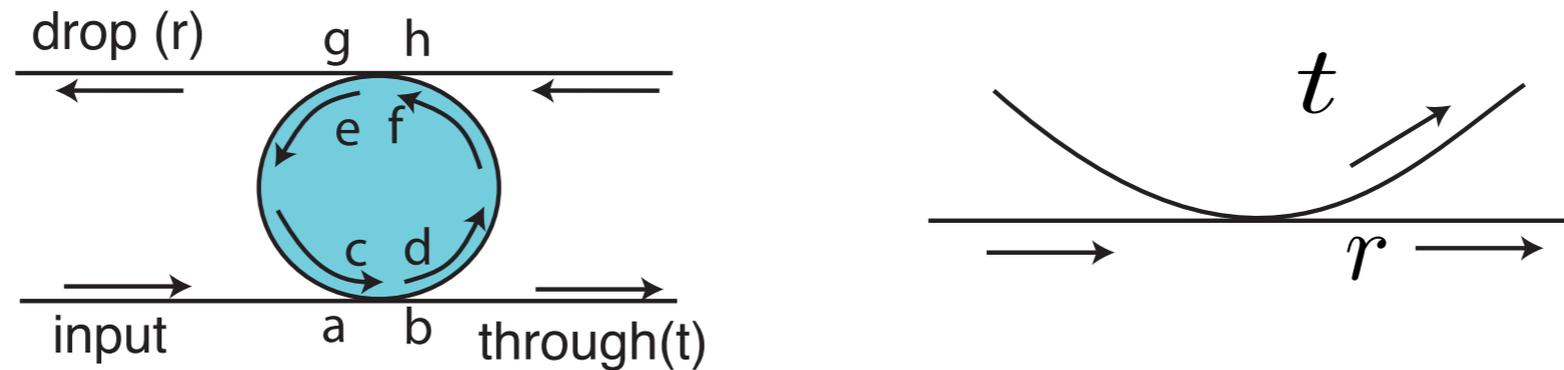
ω is the detuning

$$\mathcal{E}^{out} = \sqrt{2\kappa_{ex}}\mathcal{E}$$

$$r_{SM} = \frac{\sqrt{2\kappa_{ex}}\mathcal{E}}{\mathcal{E}_{in}e^{-i\omega t}} = \frac{2\kappa_{ex}}{i\omega - 2\kappa_{ex} - \kappa_{in}}$$



Multi-mode regime: transfer matrix formalism



$$M_{coupl} = \frac{1}{t} \begin{pmatrix} -r^2 + t^2 & r \\ -r & 1 \end{pmatrix}, \quad \begin{pmatrix} d \\ c \end{pmatrix} = M_{coupl} \begin{pmatrix} a \\ b \end{pmatrix}, \quad \begin{pmatrix} g \\ h \end{pmatrix} = M_{coupl} \begin{pmatrix} f \\ e \end{pmatrix}$$

$$|t|^2 + |r|^2 = 1$$

$$M_{prop} = \begin{pmatrix} e^{i\beta L/2 - \alpha' L/2} & 0 \\ 0 & e^{-i\beta L/2 + \alpha' L/2} \end{pmatrix}, \quad \begin{pmatrix} f \\ e \end{pmatrix} = M_{prop} \begin{pmatrix} d \\ c \end{pmatrix}$$

L is the total length of the resonator.

wave number: $\beta = n\omega/c$ α' : propagation loss

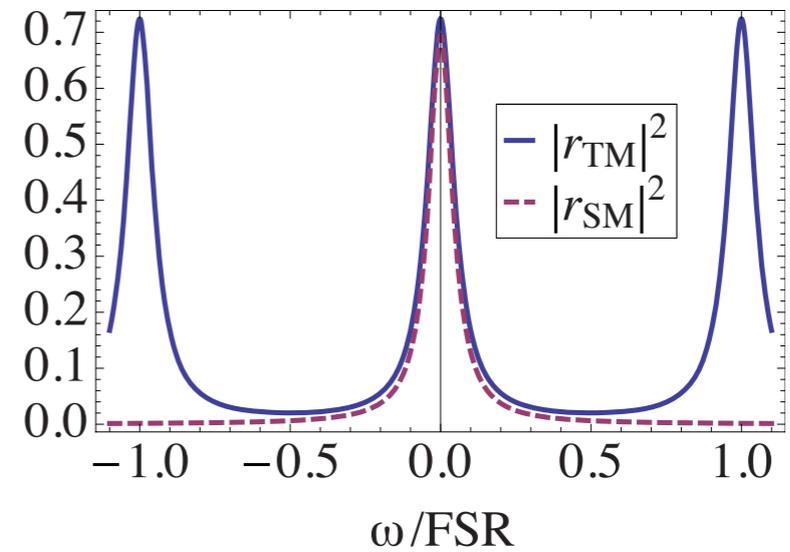
Given $a=1$ and $h=0$, we can find (b, g)

Assume the coupling is weak: $r \rightarrow \sqrt{1 - \epsilon^2}, t \rightarrow i\epsilon$

$$\epsilon \ll 1$$

$$FSR = 2\pi c/L$$

$$r_{TM} = \frac{\epsilon^2}{i\beta L - \alpha' L - \epsilon^2}$$

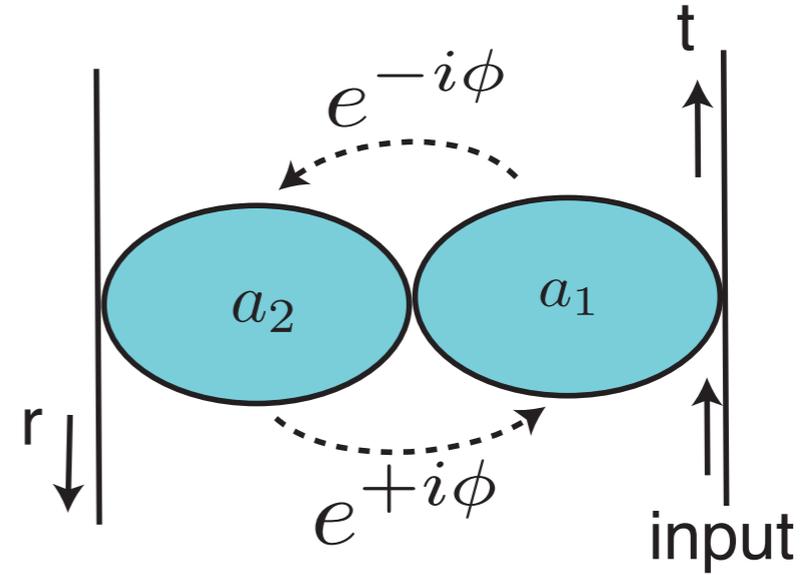


$$r_{SM} = \frac{\sqrt{2\kappa_{ex}}\mathcal{E}}{\mathcal{E}_{in}e^{-i\omega t}} = \frac{2\kappa_{ex}}{i\omega - 2\kappa_{ex} - \kappa_{in}}$$

$$\epsilon^2 \rightarrow \frac{4\pi\kappa_{ex}}{FSR}, \quad \alpha' L \rightarrow \frac{2\pi\kappa_{in}}{FSR}, \quad \beta L \rightarrow 2\pi \frac{\omega}{FSR}$$

Desired Hamiltonian, without coupling waveguides:

$$H = -J\hat{a}_2^\dagger\hat{a}_1e^{-i\phi} - J\hat{a}_1^\dagger\hat{a}_2e^{+i\phi}$$

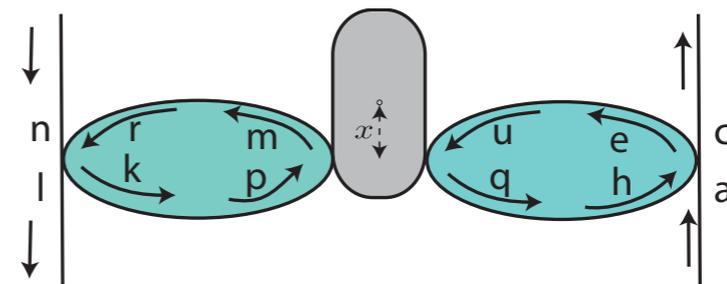


With coupling waveguides:

$$\frac{d}{dt} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \begin{pmatrix} -\kappa_{in} - \kappa_{ex} & iJe^{+I\phi} \\ iJe^{-I\phi} & -\kappa_{in} - \kappa_{ex} \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} - \sqrt{2\kappa_{ex}} \begin{pmatrix} \mathcal{E}_{in} \\ 0 \end{pmatrix}$$

$$a_2^{out} = \sqrt{2\kappa_{ex}}a_2, a_1^{out} = \mathcal{E}_{in} + \sqrt{2\kappa_{ex}}a_1$$

$$r_{SM} = \frac{\sqrt{2\kappa_{ex}}a_2}{\mathcal{E}_{in}} = -\frac{2ie^{-i\phi}J\kappa_{ex}}{J^2 + (i\omega - \kappa_{ex} - \kappa_{in})^2}$$

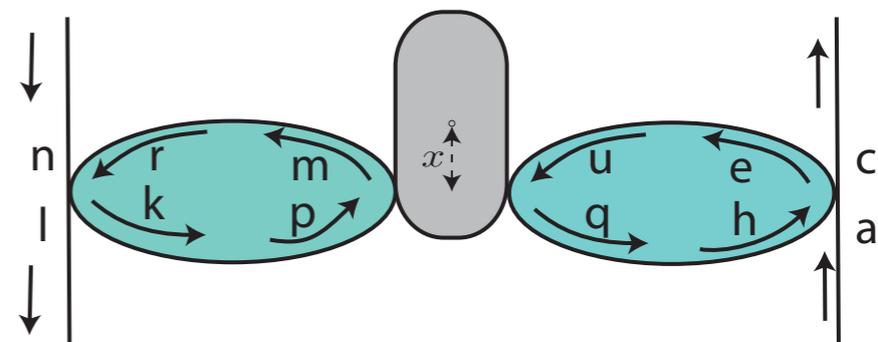


From transfer matrix approach: $r_{TM} = \frac{2e^{-i\phi}J\kappa_{ex}}{2J(\kappa_{ex} + \kappa_{in} - i\omega) \cos(\beta\eta) - i(J^2 + (\kappa_{ex} + \kappa_{in} - i\omega)^2) \sin(\beta\eta)}$

$$\epsilon_1^2 \rightarrow \frac{4\pi\kappa_{ex}}{\text{FSR}}, \quad \epsilon_2^2 \rightarrow \frac{4\pi J}{\text{FSR}}, \quad \alpha'L \rightarrow \frac{2\pi\kappa_{in}}{\text{FSR}}, \quad \beta L \rightarrow 2\pi \frac{\omega}{\text{FSR}}, \quad \beta x \rightarrow \phi$$

$$r_{SM} = \frac{\sqrt{2\kappa_{ex}}a_2}{\mathcal{E}_{in}} = -\frac{2ie^{-i\phi}J\kappa_{ex}}{J^2 + (i\omega - \kappa_{ex} - \kappa_{in})^2}$$

$$r_{TM} = \frac{2e^{-i\phi}J\kappa_{ex}}{2J(\kappa_{ex} + \kappa_{in} - i\omega)\cos(\beta\eta) - i(J^2 + (\kappa_{ex} + \kappa_{in} - i\omega)^2)\sin(\beta\eta)}$$



For: $\beta\eta = \pi/2, 3\pi/2, \dots$ both expressions become equal

$$H = -J\hat{a}_2^\dagger\hat{a}_1e^{-i\phi} - J\hat{a}_1^\dagger\hat{a}_2e^{+i\phi}$$

In general, $J_{eff} \rightarrow J/\sin(\beta\eta)$

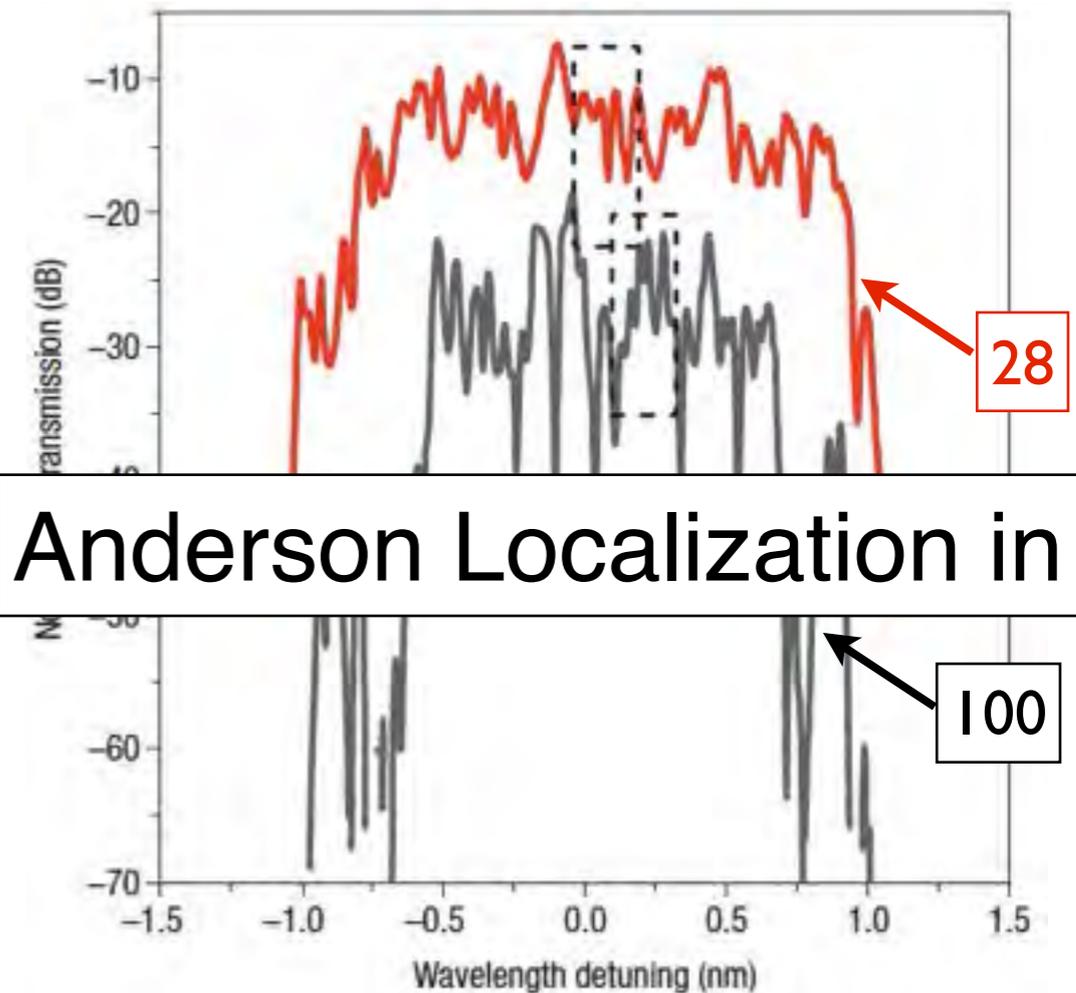
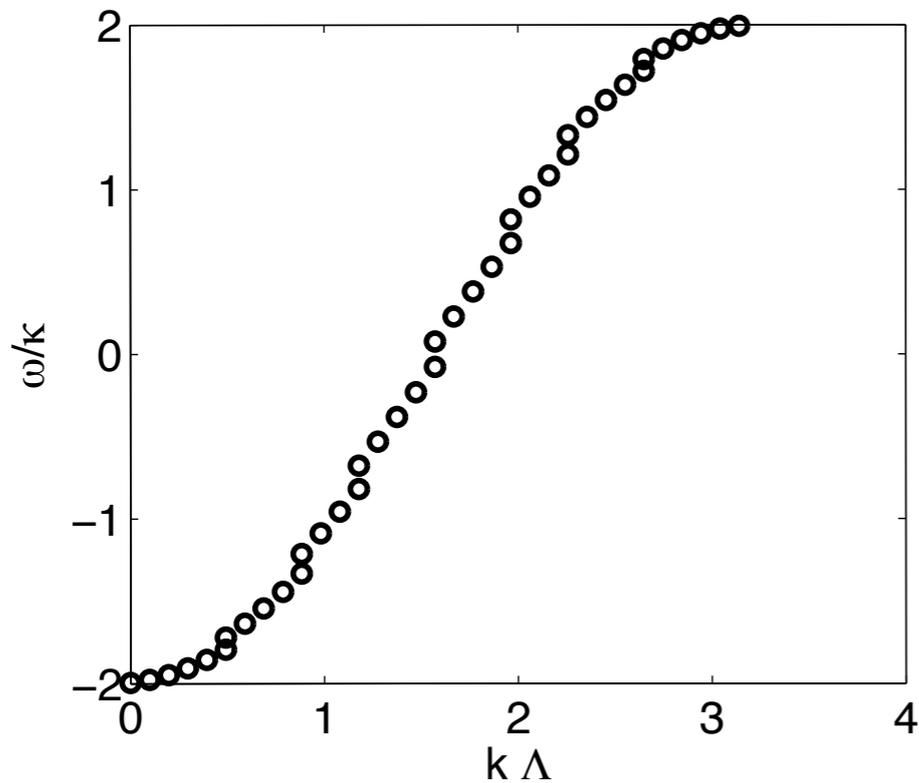
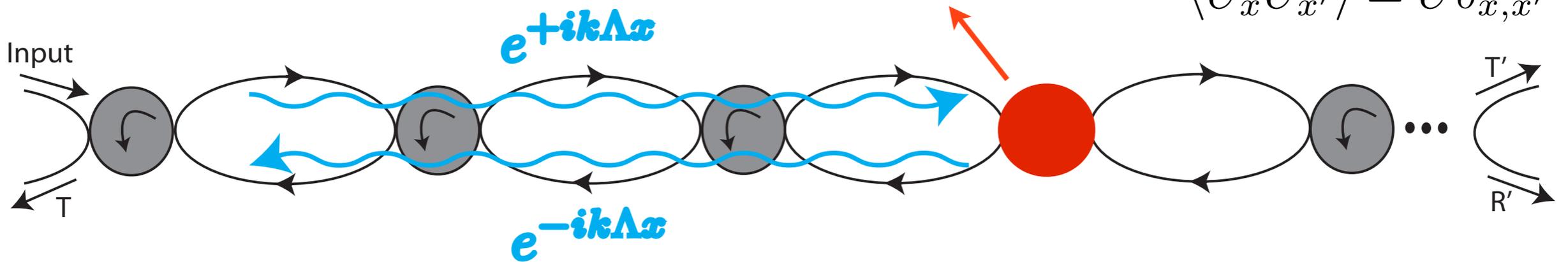
$$\omega \rightarrow \omega - J\cot(\beta\eta)$$

Coupled Resonator Optical Waveguides (CROW)

Tight-binding form:

$$H_{eff} = -\kappa \hat{a}_{x+1}^\dagger \hat{a}_x - \kappa \hat{a}_x^\dagger \hat{a}_{x+1} + U_x \hat{a}_x^\dagger \hat{a}_x$$

κ : resonator decay rate into the waveguide
 $\langle U_x \rangle = U_{x,x'}$



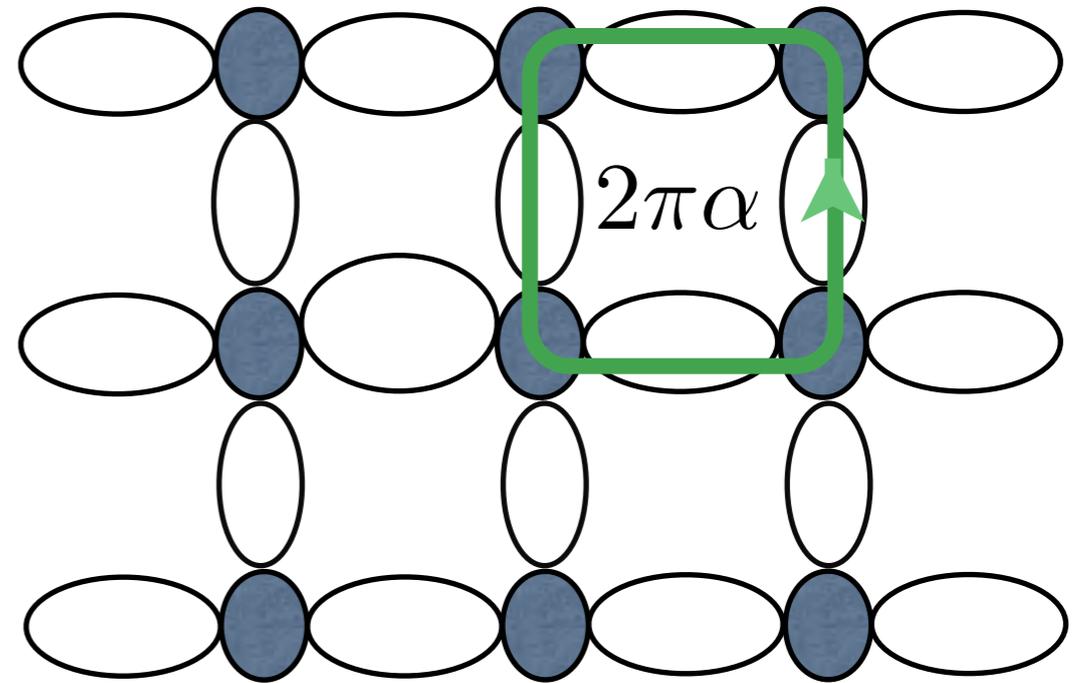
Anderson Localization in 1D

Synthetic Magnetic Field

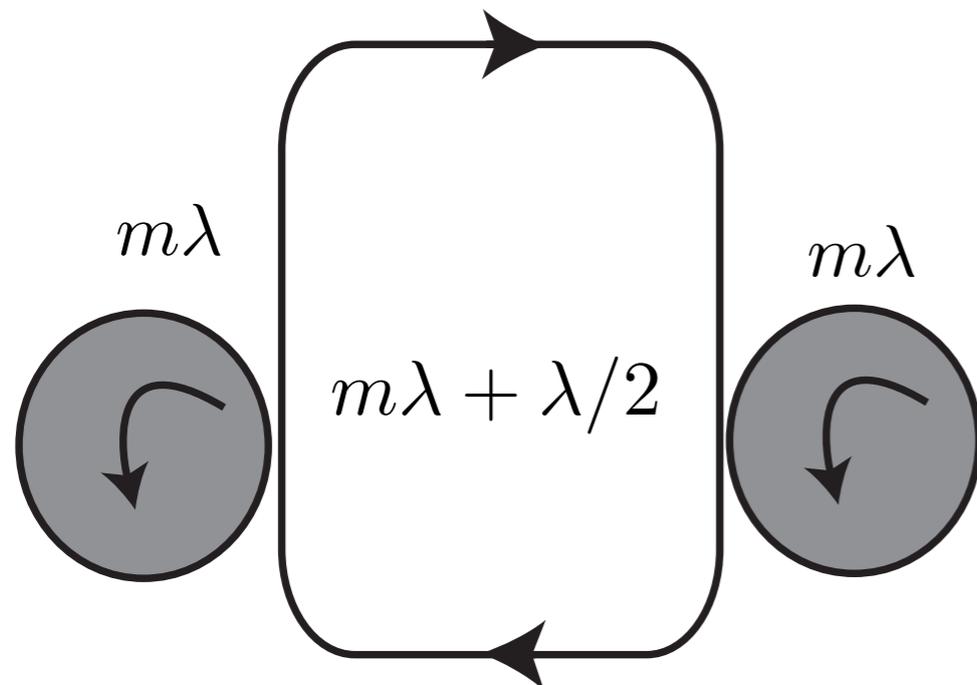
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- Tight-binding form
- Magnetic phase



Two resonator case:

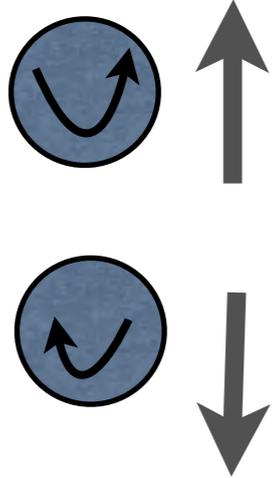


$$H_{eff} = -\kappa \hat{a}_{x+1}^\dagger \hat{a}_x e^{-2\pi i \alpha} - \kappa \hat{a}_x^\dagger \hat{a}_{x+1} e^{2\pi i \alpha}$$

MH, Demler, Lukin, Taylor
 Nat. Phys. 7, 907 (2011),
 see also Umucalilar and
 Carusotto PRA (2011)

Quantum Spin Hall

Pseudo Spin



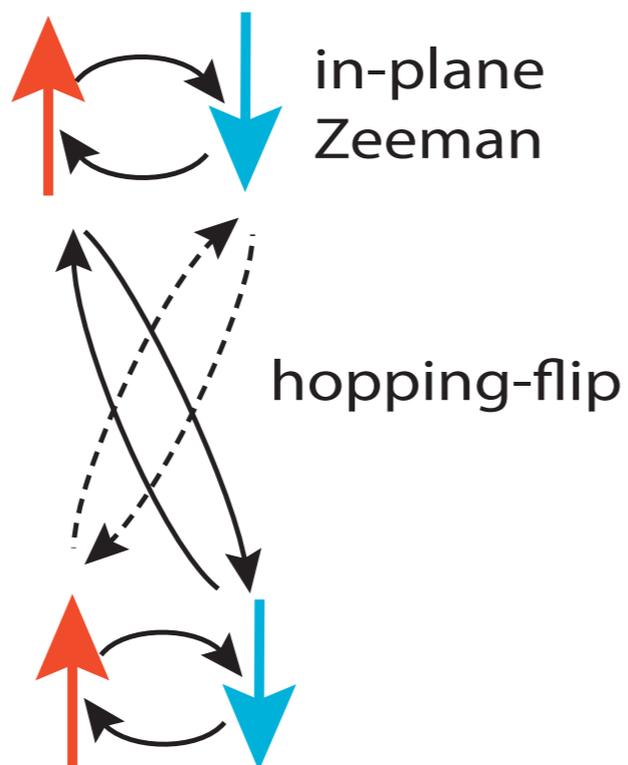
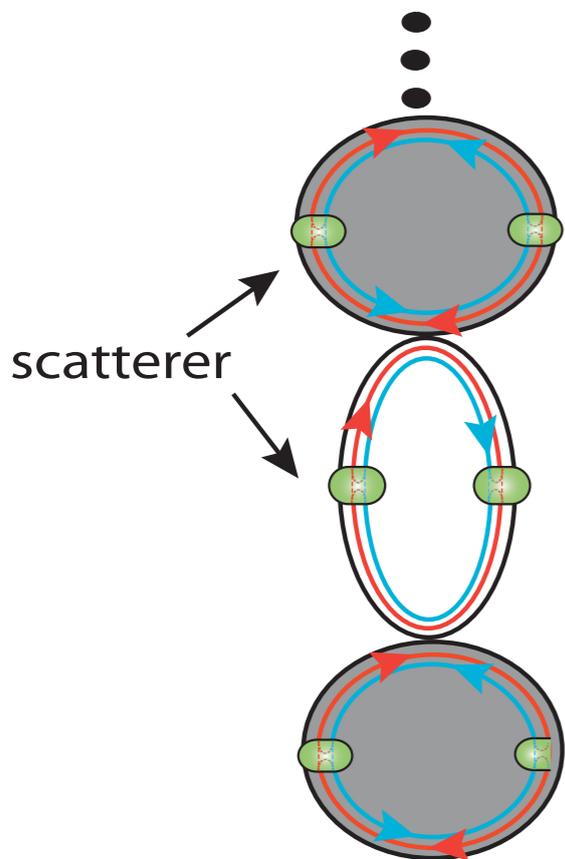
Magnetic field

$$+2\pi\alpha$$

$$-2\pi\alpha$$

$$H_{SO} \rightarrow S.L \quad \text{time-reversal invariance}$$

Adding spin-flip terms:



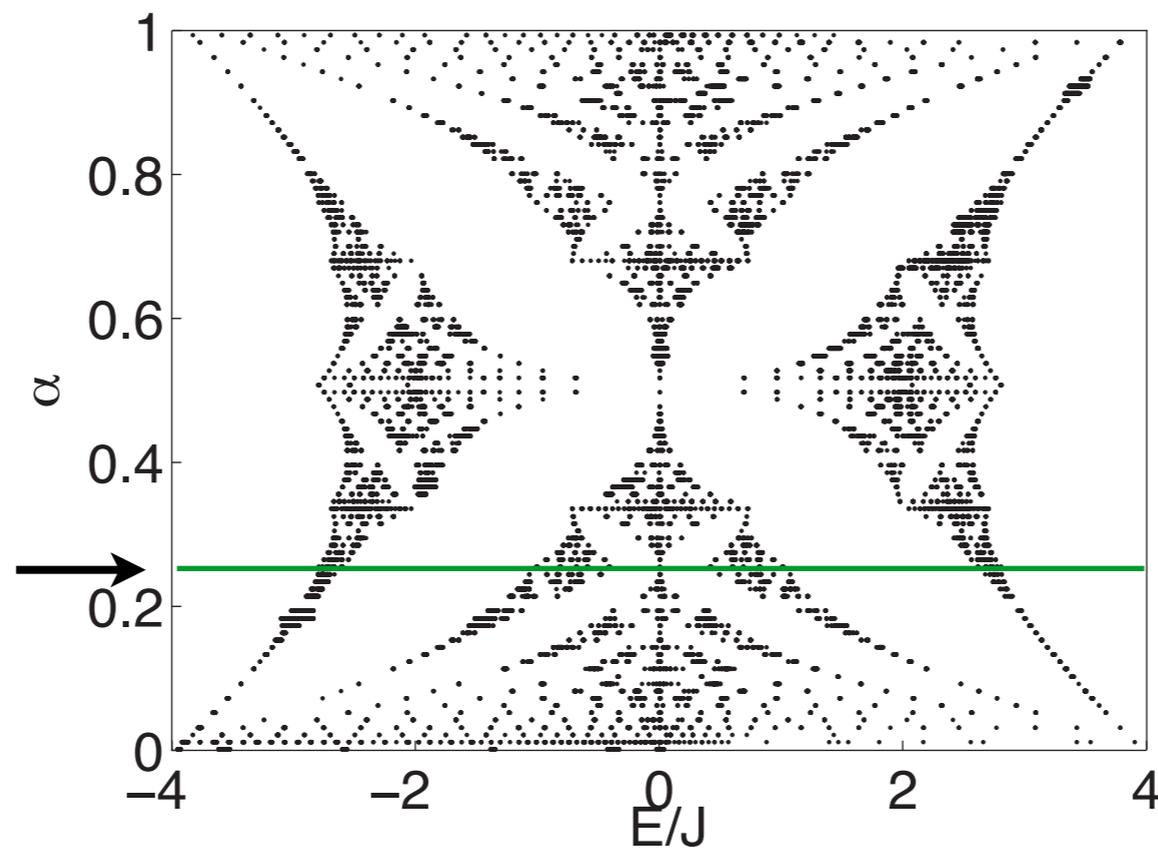
$$H = - \sum_{x,y} \begin{pmatrix} a_{\uparrow x,y+1}^\dagger & a_{\downarrow x,y+1}^\dagger \end{pmatrix} \begin{pmatrix} \kappa & 0 \\ 0 & \kappa \end{pmatrix} \begin{pmatrix} a_{\uparrow x,y} \\ a_{\downarrow x,y} \end{pmatrix} + h.c.$$

$$+ \frac{2}{\pi} \sum_{x,y} \begin{pmatrix} a_{\uparrow x,y}^\dagger & a_{\downarrow x,y}^\dagger \end{pmatrix} \begin{pmatrix} 0 & \epsilon\kappa\mathcal{F} \\ \epsilon\kappa\mathcal{F} & 0 \end{pmatrix} \begin{pmatrix} a_{\uparrow x,y} \\ a_{\downarrow x,y} \end{pmatrix}$$

$$H = - \sum_{x,y} \begin{pmatrix} a_{\uparrow x,y+1}^\dagger & a_{\downarrow x,y+1}^\dagger \end{pmatrix} \begin{pmatrix} \kappa & \kappa\epsilon \\ \kappa\epsilon & \kappa \end{pmatrix} \begin{pmatrix} a_{\uparrow x,y} \\ a_{\downarrow x,y} \end{pmatrix} + h.c.$$

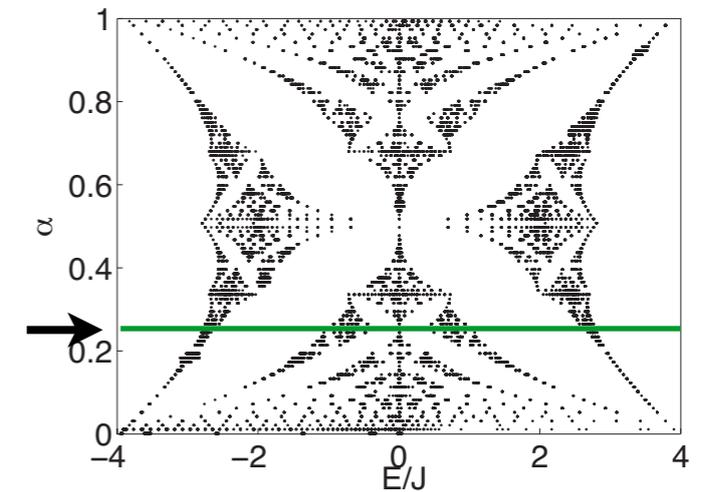
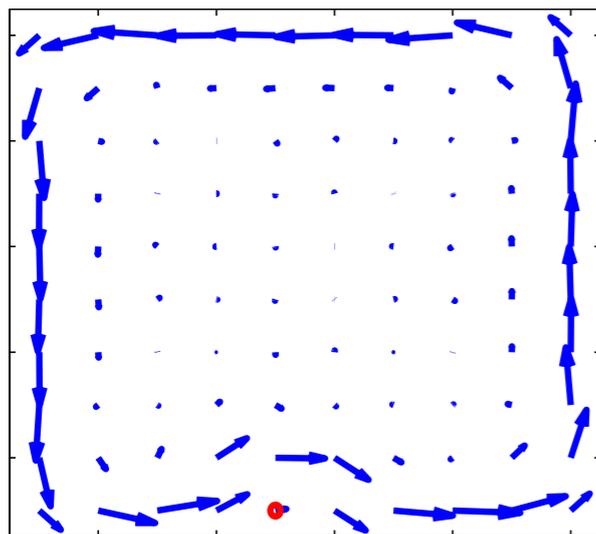
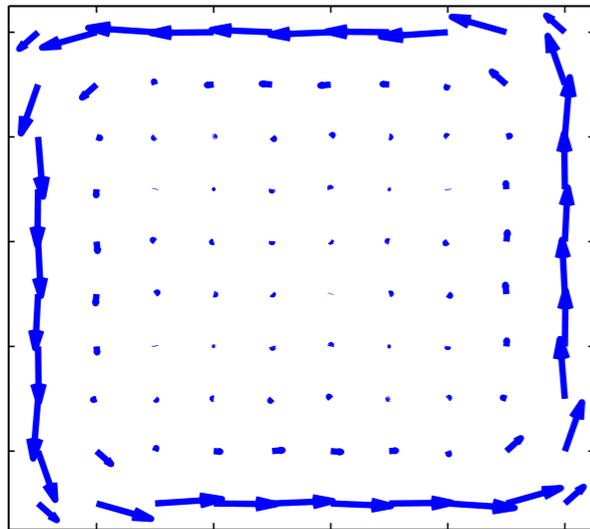
Probing photonic system

- Probing allows us to perform a “spectroscopy” of photonic states
- Energy spectrum of an infinite tight-binding with uniform magnetic field is [Hofstadter butterfly](#)
- Torus boundary condition (10x10) can simulate an infinite system



Edge states

- Finite system has edges -> New states beside Hofstadter states
- Quasi-one dimensional states that are **localized at the perimeter** of the system
- Carry a **chiral current**

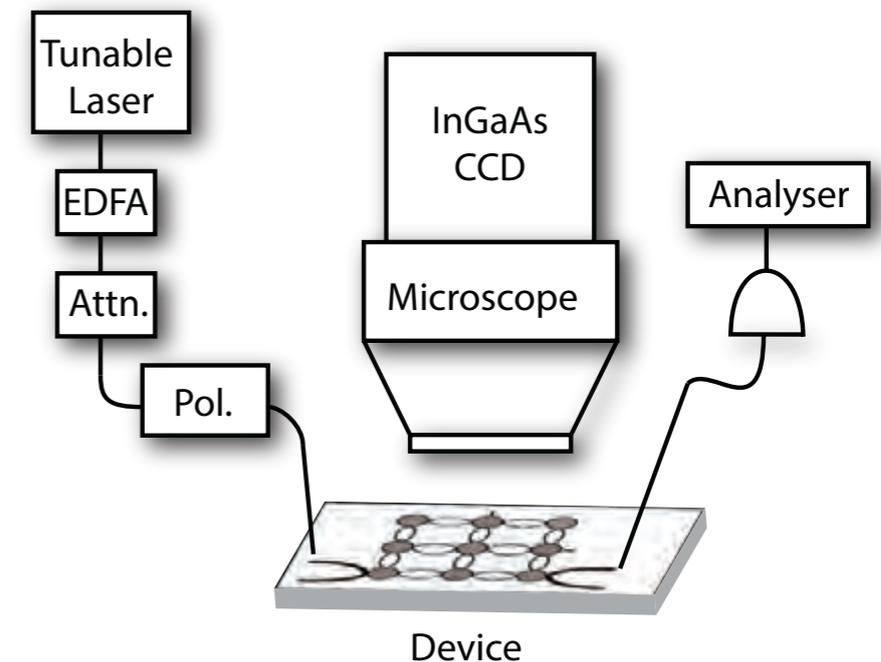
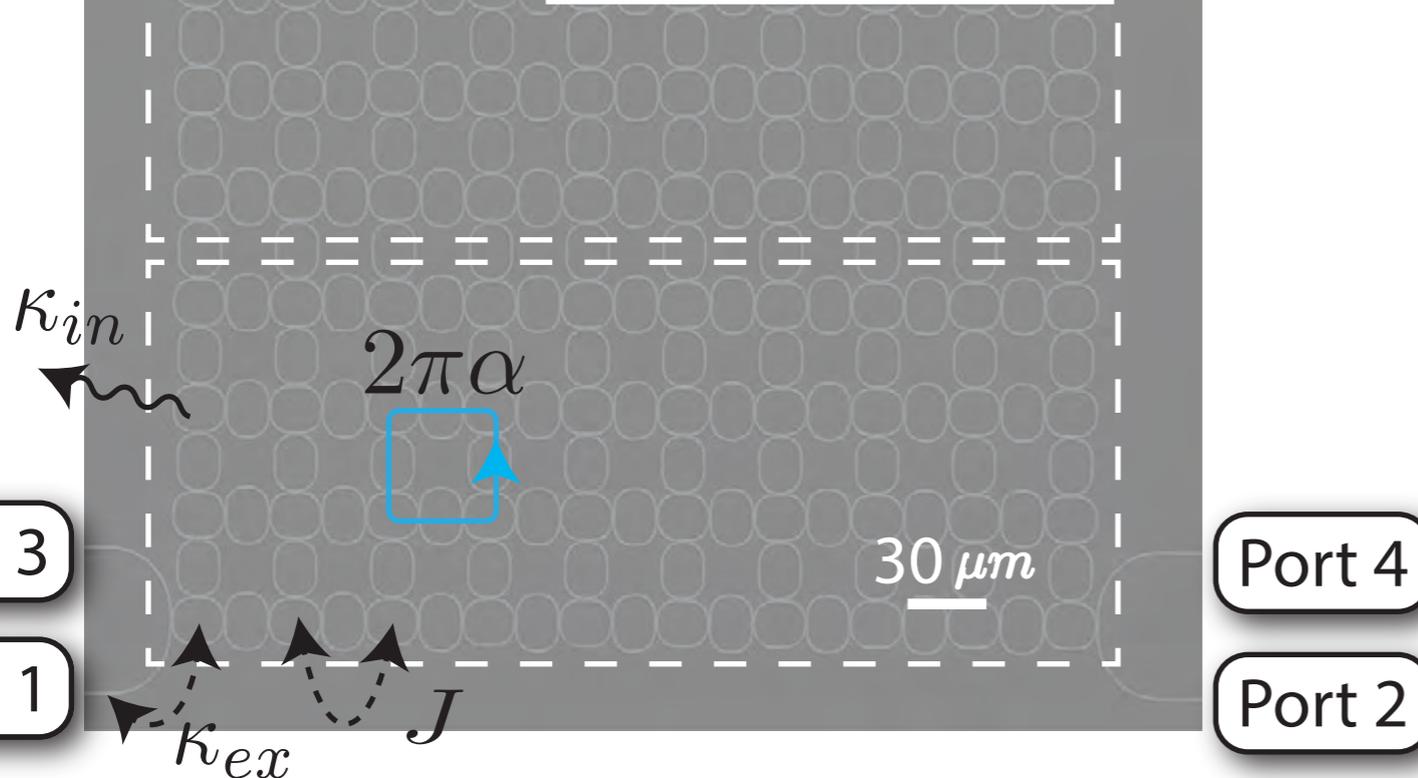
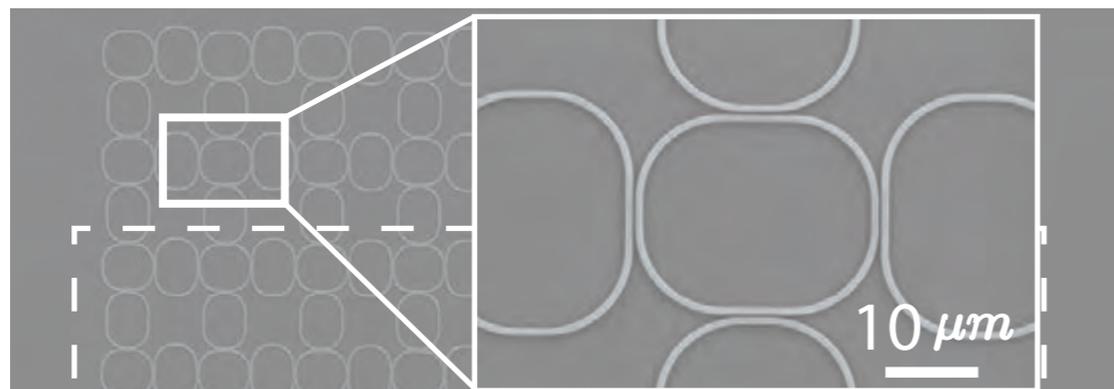
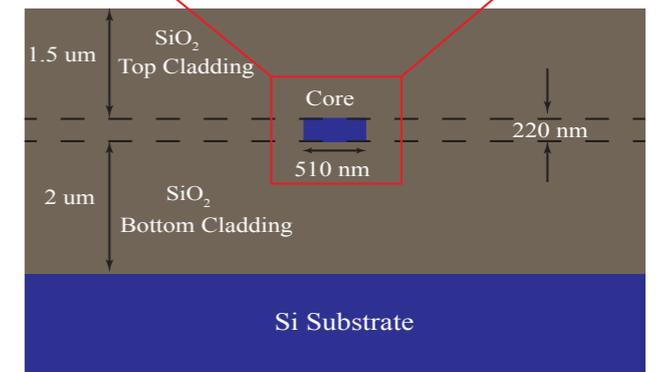
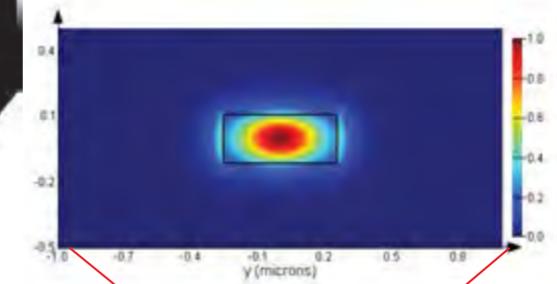


- **Robust** against non-magnetic impurities $U a_i^\dagger a_i$
- Current “gets around” the impurity rather than getting scattered in the opposite direction

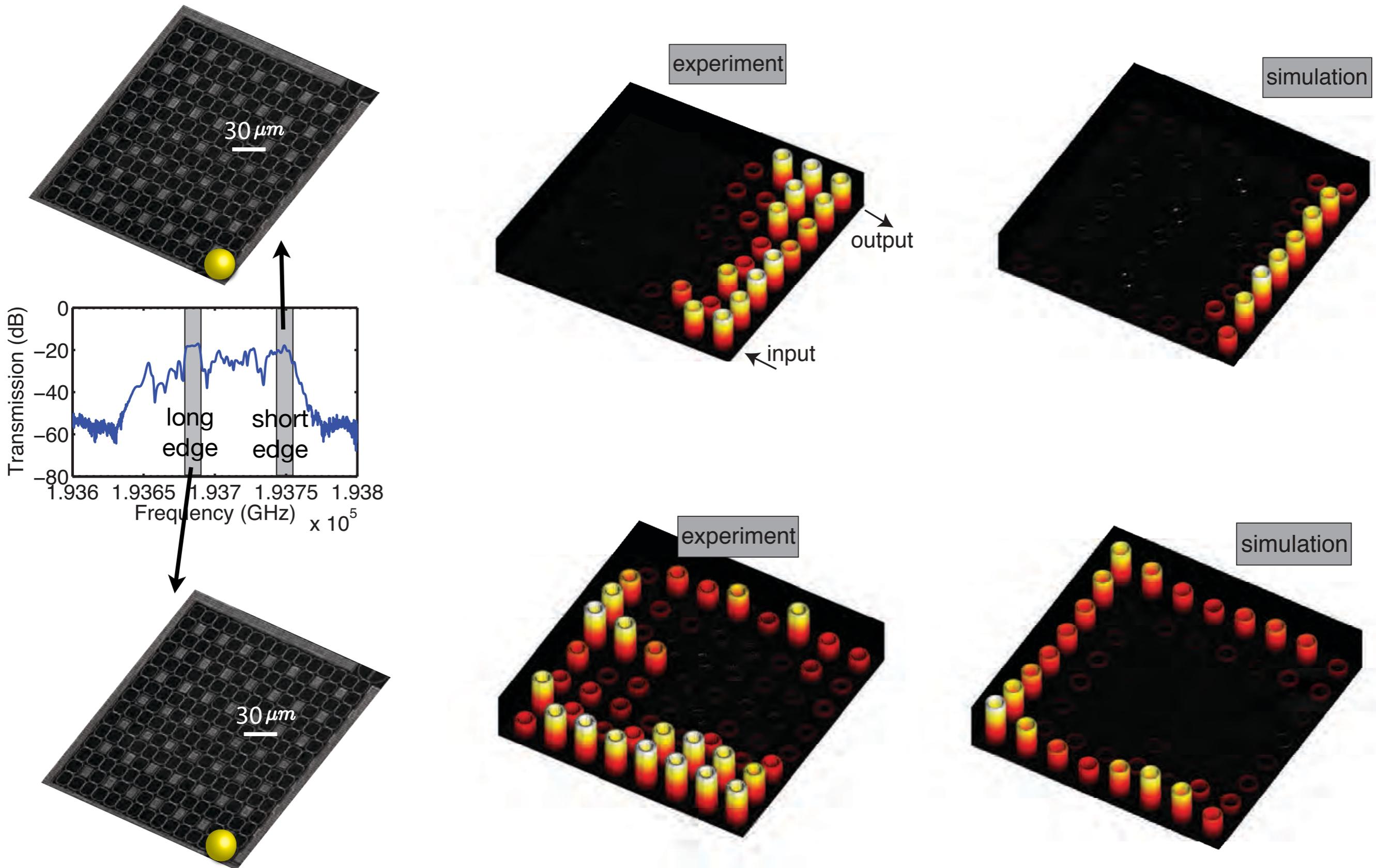
- Synthetic gauge field, formalism
- Experiment in silicon photonics
- Interacting regime: Fractional quantum Hall
 - Laughlin, Pfaffian etc.
 - Preparation
- Measuring topological invariants

Experimental realization of the gauge field

- Silicon-on-Insulator technology

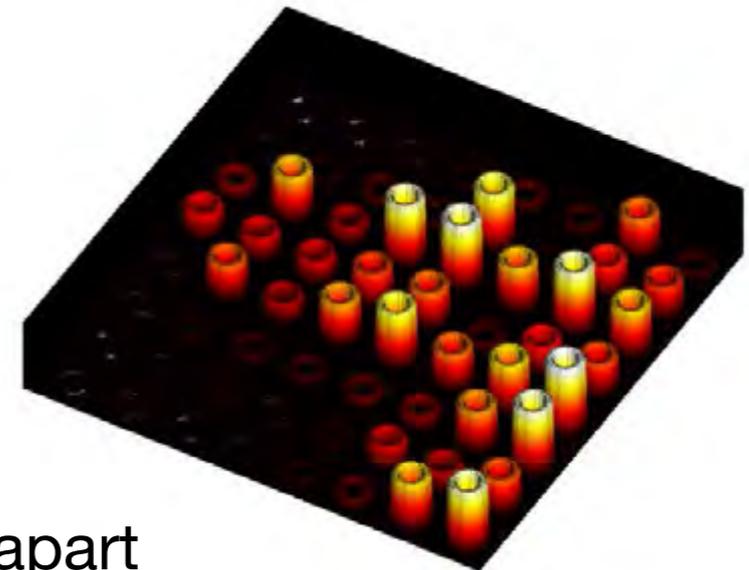
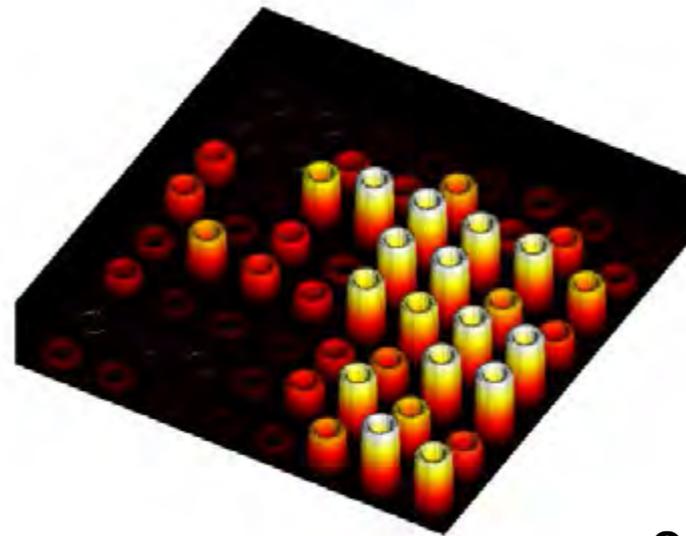


Observation of topological edge states



Topological robustness of edge state

Bulk state profile is extremely **sensitive** to frequency change.

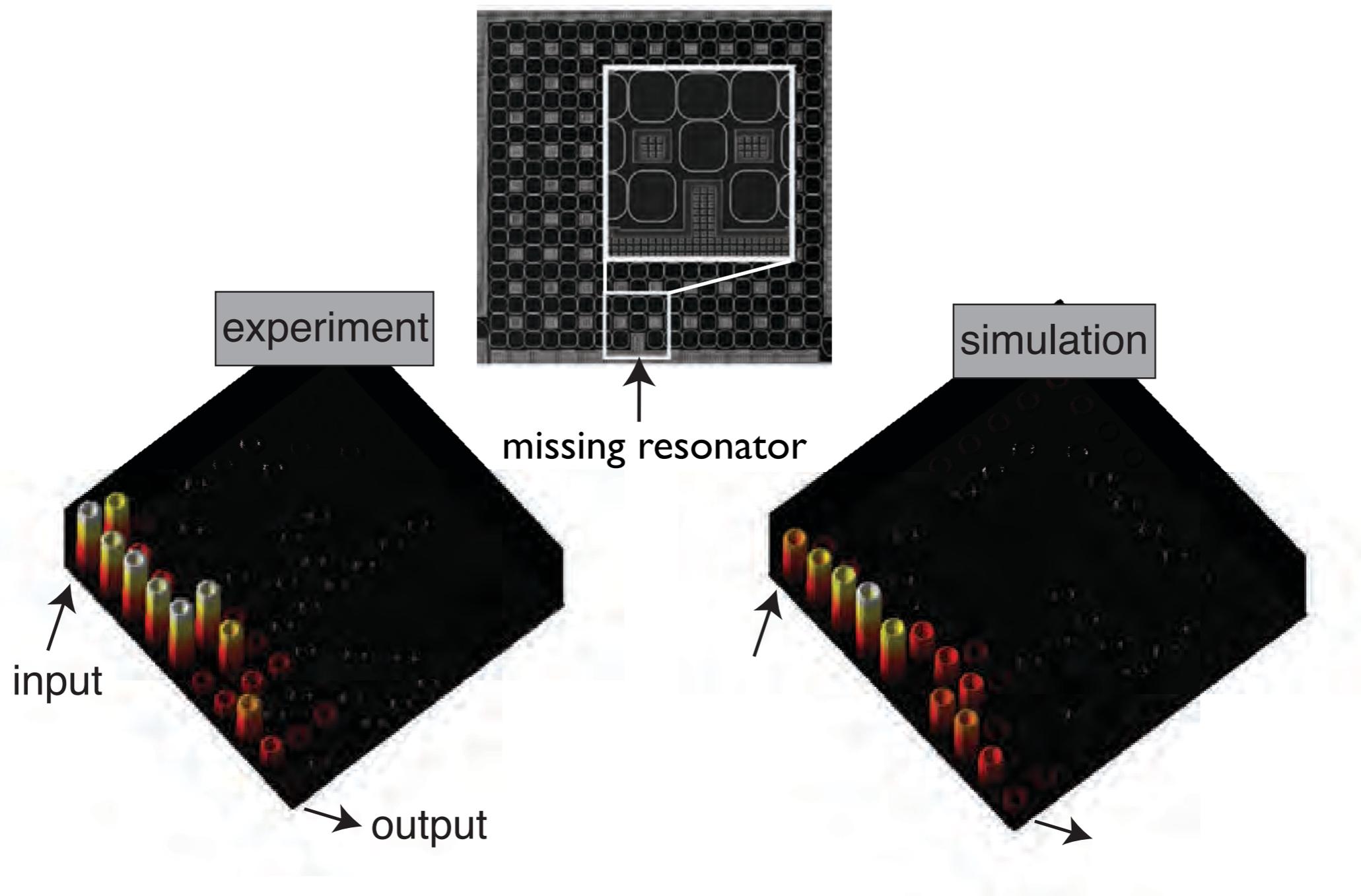


0.1GHz apart

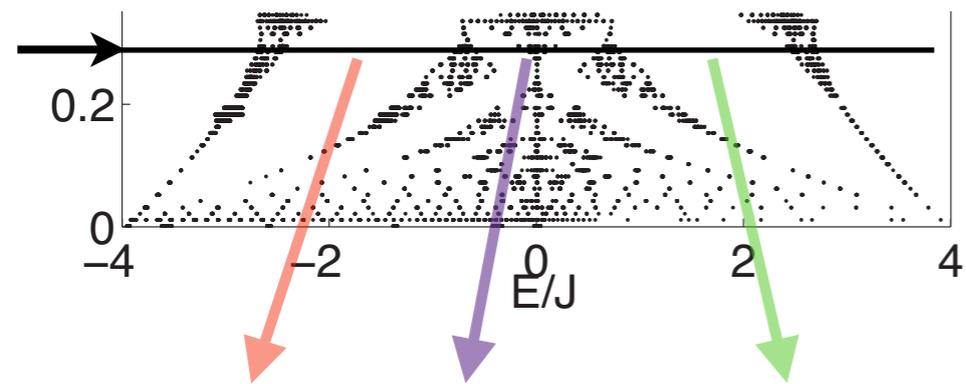
- Edge states are observed over a **broad band** (15 GHz)
- Edge state is **robust against intrinsic disorder** and the propagation profile does not change

frequency mismatch	$U = 0.8J$
loss rate inhomogeneity	$\Delta\kappa_{in} = 0.45\kappa_{in}$
coupling inhomogeneity	$\Delta J = 0.04J$
magnetic field inhomogeneity	$\Delta\alpha = 2\pi 0.08$
spin flip rate	$\beta = 0.04J$

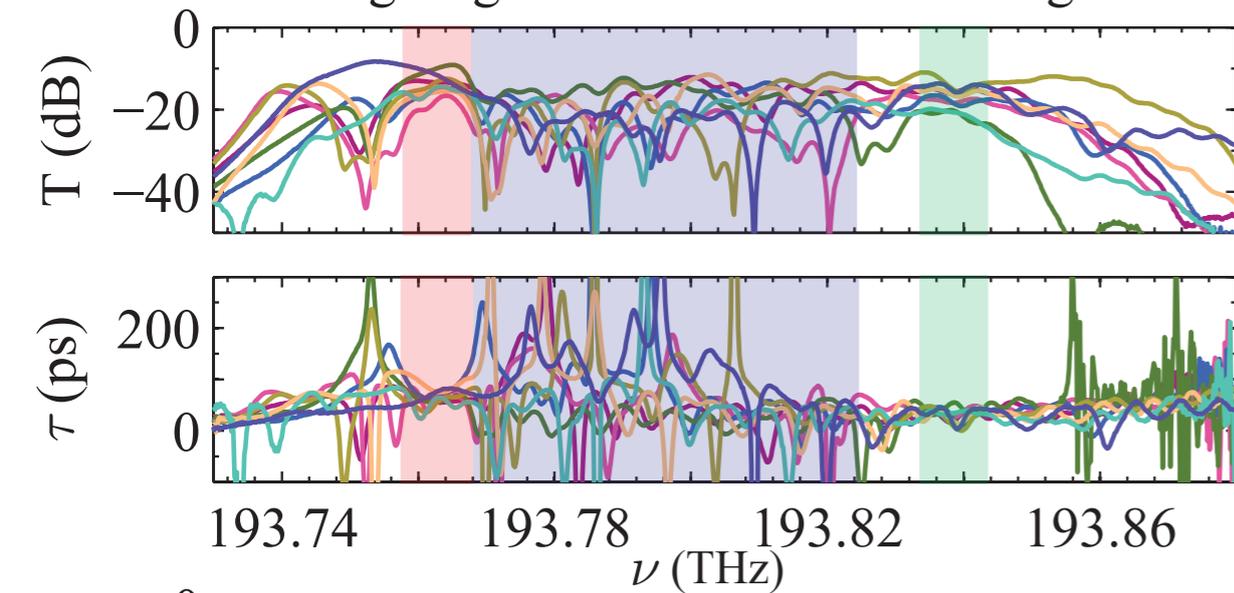
Robustness against an introduced disorder



Transport statistics

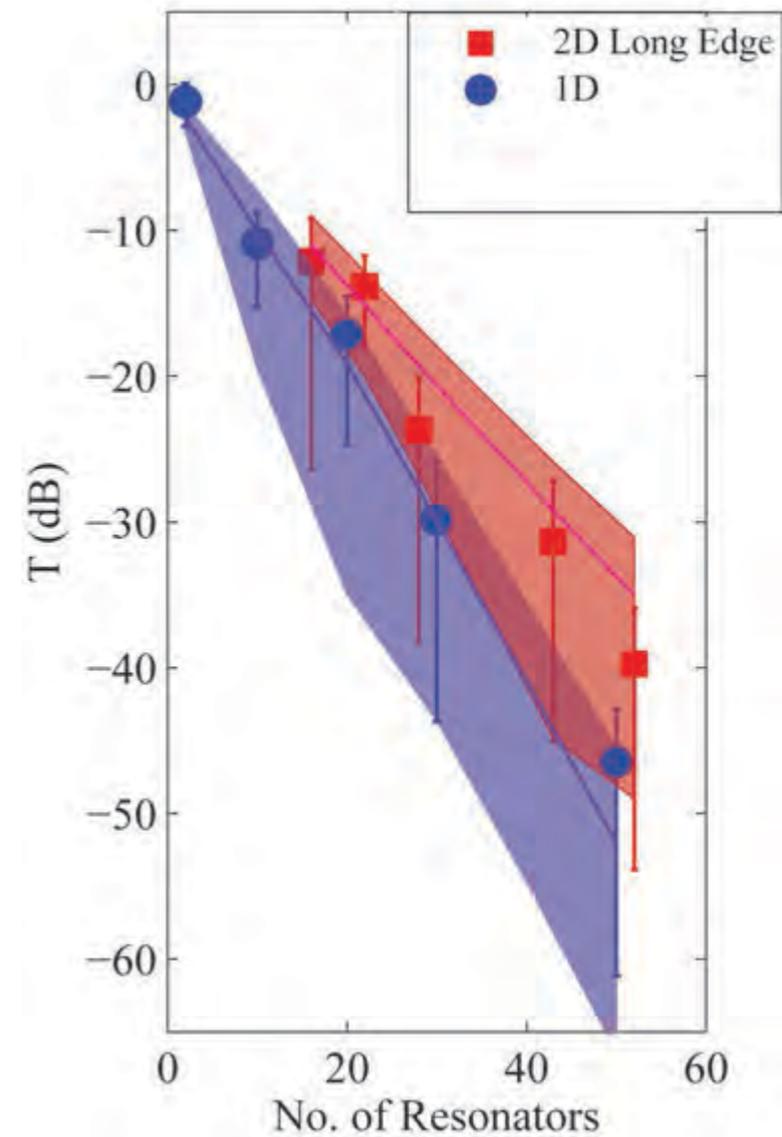


Long Edge Bulk Short Edge



15x15 arrays

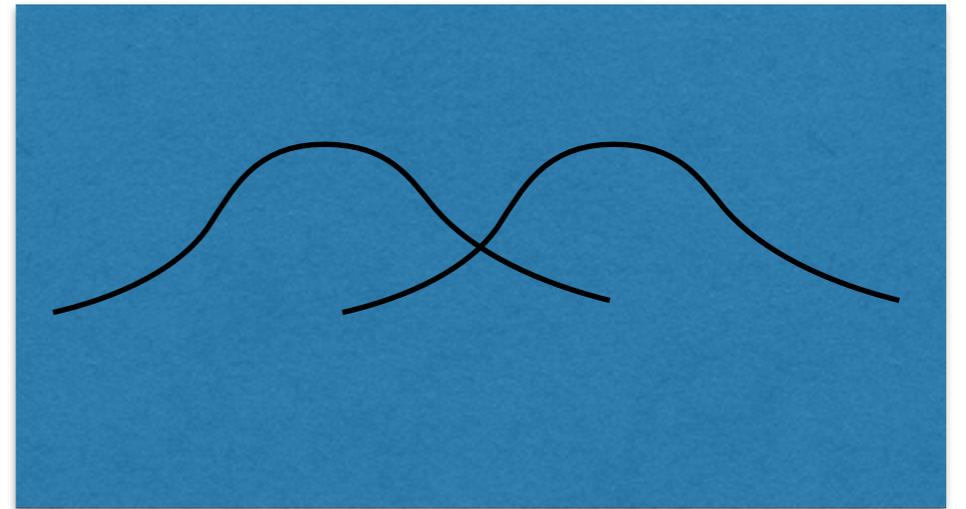
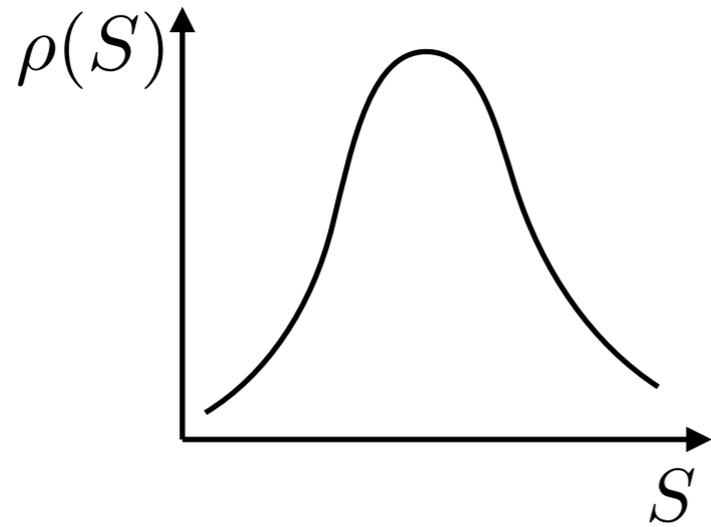
Different colors: different samples



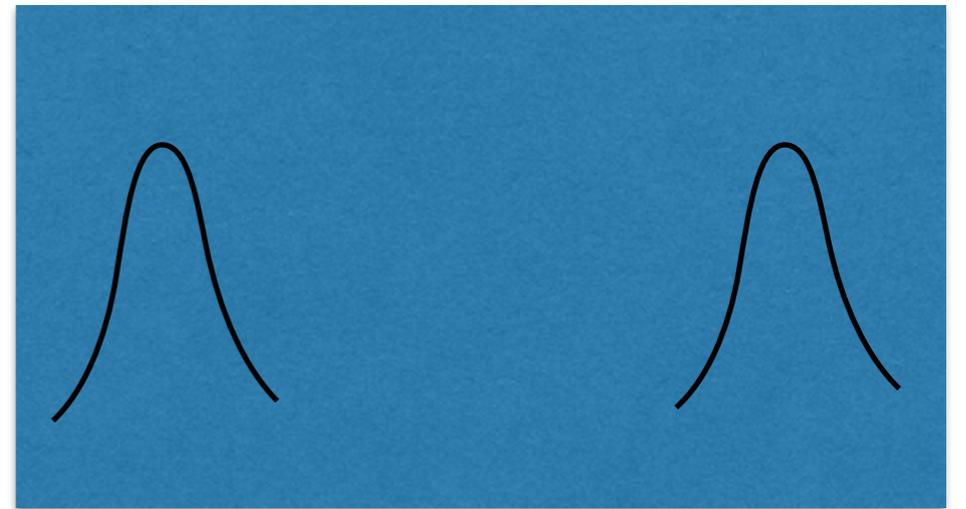
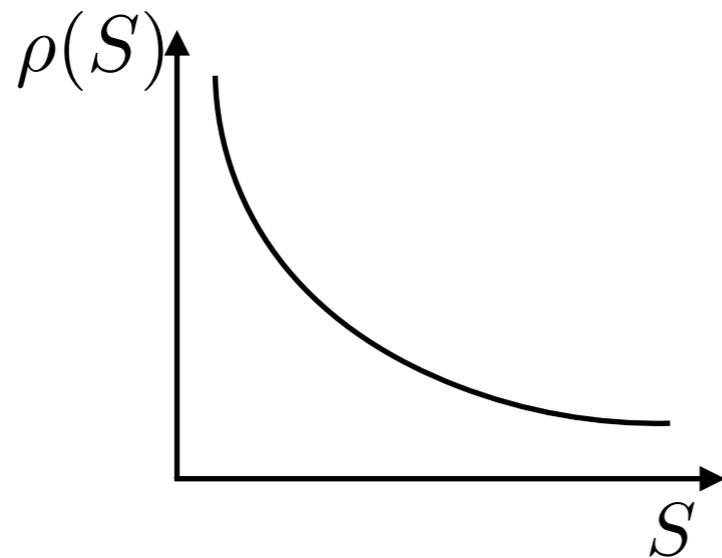
S. Mittal, J. Fan, S. Faez, A. Migdall, J. M. Taylor, and M.H.
Phys. Rev. Lett. 113, 087403 (2014)

localization and level repulsion

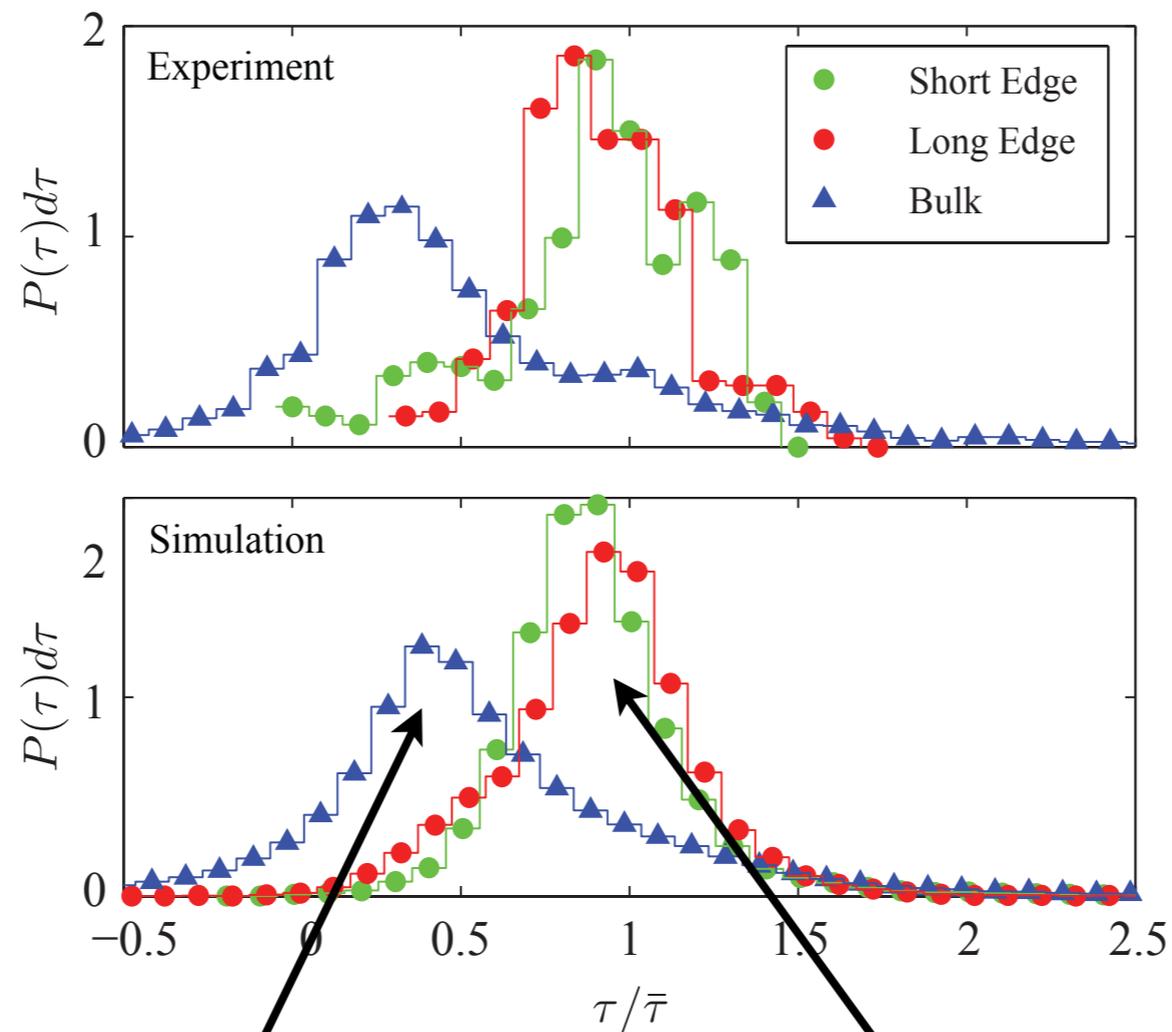
diffusive:



localized:



$$\tau(\omega) = \frac{L}{v_g(\omega)} \propto \frac{dk}{d\omega} \propto \rho(\omega)$$



localized

diffusive

Asymmetry and deviation of the peak from the mean indicates **localization of bulk states**

earlier work in 1D:

Genack Nature (2000)

Texier & Comtet (1999)

S. Mittal, J. Fan, S. Faez, A. Migdall, J. M. Taylor, and M. H.

Phys. Rev. Lett. 113, 087403 (2014)

Recent publications exploring topological properties of light

- YE Kraus, Y Lahini, Z Ringel, M Verbin, O Zilberberg - *Physical Review Letters*, 2012
- L. Lu, L. Fu, J. Joannopoulos and M. Soljacic *Nature Photonics* 7, 294–299 (2013)
- V Yannopoulos *New Journal of Physics* (2012)
- K Fang, Z Yu, S Fan - *Nature Photonics* (2012)
- M Rechtsman, et al. - *Nature Photonics* (2012)
- A. Khanikaev, S. Mousavi, W. Tse, M. Kargarian, A. MacDonald, G. Shvets, *Nature Material* (2012)
- W. Yang, Z. Yin, Z. Chen, S. Kou, M. Feng, C. Oh - *Physical Review A* (2012)
- M. Verbin, O. Zilberberg, Y. Kraus, Y. Lahini, and Y. Silberberg *Physical Review Letter* (2013)
- MC Rechtsman, et al. - *Nature* (2013)
- G. Jiang, Y. Chong *Physical Review Letters* (2013)
- V. G. Sala, D. D. Solnyshkov₃, T. Jacqmin₁, A. Lemaître₁, H. Terças₃, M. Abbarchi, E. Galopin, I. Sagnes, J. Bloch, G. Malpuech, A. Amo, *Arxiv* (2014)
- Jia Ningyuan, Ariel Sommer, David Schuster, Jonathan Simon *Arxiv* (2013)
- L. Tzuang ... M. Lipson - *Nature Photonics* (2014)

